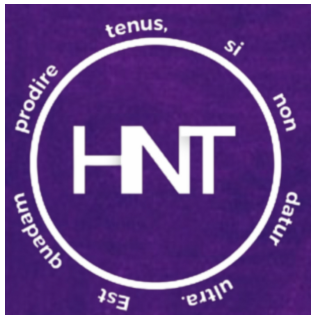


# Frobenious structures and restricred (2+1)-TQFTs

Dušan Đorđević, University of Belgrade



Фонд за науку  
Републике Србије



# In collaboration with:



Zoran Petrić



Danica Kosanović



Jovana Nikolić

- Studying QFT is hard: perturbative methods, renormalization, infinite number of degrees of freedom. Usual guide: **symmetries!**



- There are cases in condensed matter physics when certain properties of a gapped system depend only on topology!
- We can ask for a nontrivial field theory that still works with finite dimensional Hilbert spaces.

⇒ TQFT

- One physical example: Chern Simons! Take example of SU(2) gauge group and  $A$  is a gauge connection, that is locally a Lie algebra valued one form.

$$S = \int \text{Tr} (A \wedge dA + A \wedge A \wedge A) \quad \Rightarrow \quad F = 0$$

- Is this theory trivial? No! It is used in condensed matter to describe QHE, in knot theory to compute Jones polynomials and in 3D gravity.
- Wilson loop observable  $\rightarrow$  can detect nontrivial topology (anyons)!



$$\text{Tr}_R \underbrace{P e^{\oint A}}_{\text{holonomy}}$$



- Previous consideration was classical, can we go to **quantum** level?
- We will use a formal mathematical approach, first introduced by Atiyah [Atiyah '89].
- A TQFT is a (strong) symmetrical monoidal functor from the category of cobordisms to the category of vector spaces.
- This approach has an advantage that is well accessible to pure mathematics; it does not require one to deal with path integrals.
- But, what does this means?

$$\int \mathcal{D}A_\mu e^{-S(A_\mu)}$$

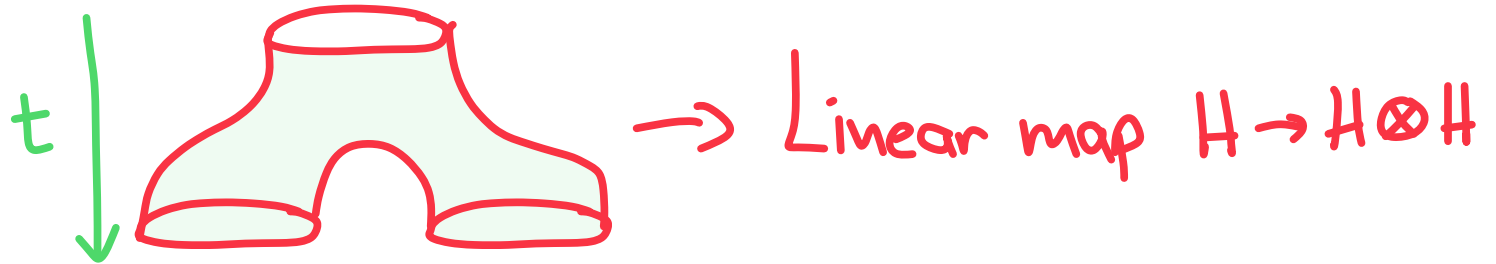
**Def:** A category consist of objects  $(a, b, c, \dots)$  and arrows  $(f, g, h, \dots)$ , together with rules:

- For every arrow  $f$ , there exist objects  $\text{dom}(f)$ ,  $\text{cod}(f)$
- For arrows  $f: a \rightarrow b$  and  $g: b \rightarrow c$  there is an arrow  $g \circ f: a \rightarrow c$
- For every object  $a$  there is an arrow  $\mathbf{1}_a: a \rightarrow a$
- Composition is associative
- For every arrow  $f: a \rightarrow b$  we have  $f \circ \mathbf{1}_a = \mathbf{1}_b \circ f$
  
- Important example:  $\mathbf{Vect}_{\mathbb{C}}$ : objects are complex vector spaces, arrows are linear maps between vector spaces.
- **Functor**: mapping between categories (think of group representations).

- Let us illustrate everything in 2D:



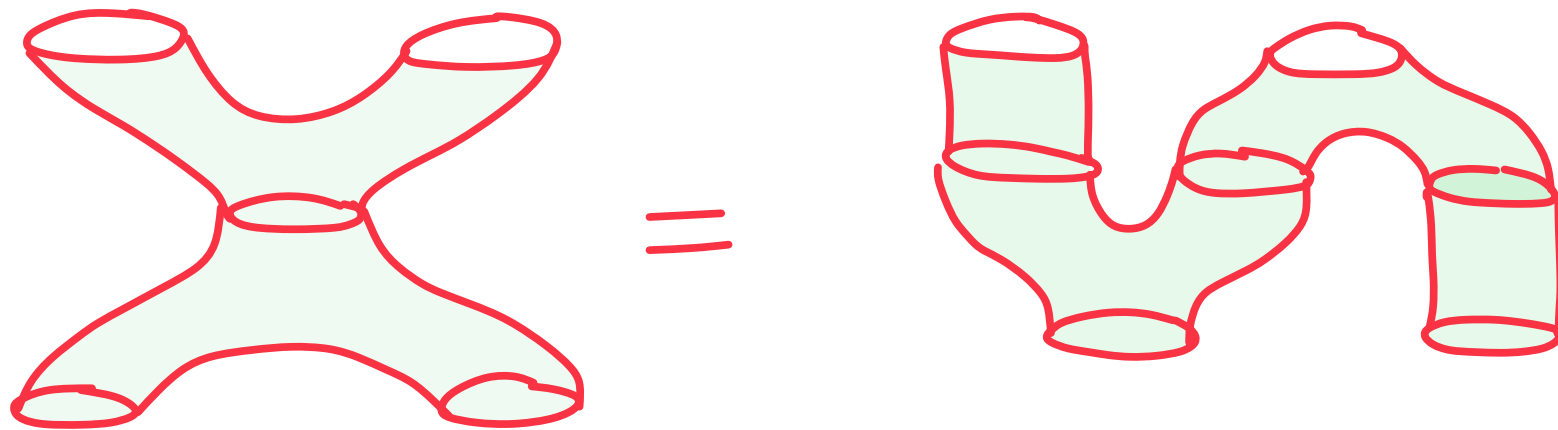
- Cobordisms!



- In two dimensions, it is well-known that we can generate the whole theory from the following cobordisms

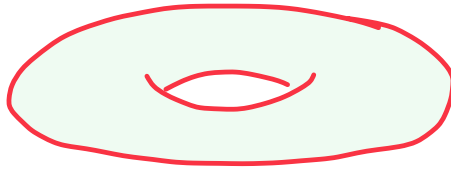


- Everything should depend only on “topology”!



Commutative Frobenius algebra

- What happens in three dimensions?
- Three manifolds are much harder to study. Also, objects of **3Cob** are closed, oriented surfaces, with much richer structure than  $S^1$ .



- The main character in our story will be torus, so let's focus on it. First, define a **mapping class group** as the group of orientation preserving homeomorphisms modulo isotopy.
- It is a well known fact that, for torus, this group is  $SL(2, \mathbb{Z})$ .

- Righthanded Dehn twists on a:



- Matrix representation

$$D_a := \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix} \quad \text{and} \quad D_b := \begin{pmatrix} 1 & 0 \\ -1 & 1 \end{pmatrix}$$

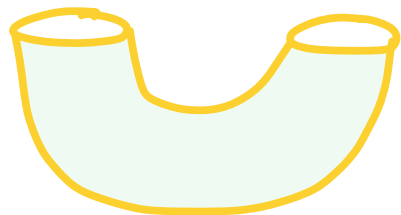
- Another useful set of generators of  $SL(2, \mathbb{Z})$ , satisfying  $S^2 = (ST)^3, S^4 = \text{Id}$

$$S = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}, \quad T = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}.$$

- Motivation of our work: **faithful TQFT**. In dimension one, almost all TQFTs are faithful [Telebaković Onić '20]. In two dimensions, there exists a faithful TQFT [Telebaković Onić, Petrić, Gajović '20]. What happens in three dimensions?
- Part of this story is well-known: **extended TQFT**. From physical perspective, they are interesting and important. Mathematically, they are based on a structure of a modular tensor category. Example: Reshetikin-Turaev (RT) TQFT.
- Unfortunately, those TQFTs are not faithful [Funar '12], there exist some torus bundles that are not distinguished by RT invariants. We therefore seek to find some representations of  $SL(2, \mathbb{Z})$  that are not connected with MTC.

- Actually, similar to the classification of 2D TQFTs in terms of Frobenius algebras, 3D TQFTs can be classified using J algebras [Juhasz '18]. However, J algebras are much harder to work with and we therefore seek to define a restricted TQFT.

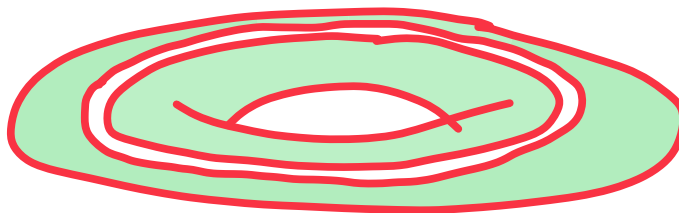
- We consider only a subcategory of **3Cob**. Today, we focus on  $\beta$ MCG.



$$\beta \circ (D_a \otimes \mathbf{1}) = \beta \circ (\mathbf{1} \otimes D_b).$$

$$\beta \circ (D_b \otimes \mathbf{1}) = \beta \circ (\mathbf{1} \otimes D_a).$$

- General hint: thick torus!





- Here we have: the identity, the pairing  $\beta$ , the copairing  $\gamma$  and the mapping cylinder.
- Examples of representations of  $SL(2, \mathbb{Z})$  that are able to distinguish between torus bundles:

$$S = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \end{pmatrix}$$

or

$$S = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}$$

or

...

$$T = \begin{pmatrix} 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \end{pmatrix}$$

$$T = \begin{pmatrix} 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}$$



- We take the following representation for MCG:

$$D_a = \begin{pmatrix} 0 & 0 & -\frac{e^{-\frac{i\pi}{6}}}{\sqrt{2}} & 0 & -\frac{1}{2}ie^{-\frac{i\pi}{6}} & 0 & 0 \\ \frac{e^{-\frac{i\pi}{6}}}{\sqrt{2}} & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & e^{-\frac{i\pi}{6}} & 0 & -i\sqrt{2}e^{-\frac{i\pi}{6}} & -\frac{e^{-\frac{i\pi}{6}}}{\sqrt{2}} \\ 0 & -\frac{1}{2}e^{-\frac{i\pi}{6}} & e^{-\frac{i\pi}{6}} & 0 & -\frac{ie^{-\frac{i\pi}{6}}}{\sqrt{2}} & 0 & 0 \\ 0 & 0 & 0 & -i\sqrt{2}e^{-\frac{i\pi}{6}} & 0 & -2e^{-\frac{i\pi}{6}} & -ie^{-\frac{i\pi}{6}} \\ 0 & -\frac{ie^{-\frac{i\pi}{6}}}{2\sqrt{2}} & -\frac{ie^{-\frac{i\pi}{6}}}{\sqrt{2}} & 0 & -\frac{1}{2}e^{-\frac{i\pi}{6}} & 0 & 0 \\ 0 & 0 & 0 & -\frac{e^{-\frac{i\pi}{6}}}{2\sqrt{2}} & 0 & -\frac{1}{2}ie^{-\frac{i\pi}{6}} & 0 \end{pmatrix} \rightarrow \text{Works!}$$

$$D_b = \begin{pmatrix} 0 & \sqrt{2}e^{-\frac{i\pi}{6}} & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -\frac{1}{2}e^{-\frac{i\pi}{6}} & 0 & -\frac{ie^{-\frac{i\pi}{6}}}{\sqrt{2}} & 0 \\ -\frac{e^{-\frac{i\pi}{6}}}{2\sqrt{2}} & 0 & 0 & e^{-\frac{i\pi}{6}} & 0 & -i\sqrt{2}e^{-\frac{i\pi}{6}} & 0 \\ 0 & 0 & e^{-\frac{i\pi}{6}} & 0 & -\frac{ie^{-\frac{i\pi}{6}}}{\sqrt{2}} & 0 & -\frac{e^{-\frac{i\pi}{6}}}{\sqrt{2}} \\ -\frac{1}{2}ie^{-\frac{i\pi}{6}} & 0 & 0 & -i\sqrt{2}e^{-\frac{i\pi}{6}} & 0 & -2e^{-\frac{i\pi}{6}} & 0 \\ 0 & 0 & -\frac{ie^{-\frac{i\pi}{6}}}{\sqrt{2}} & 0 & -\frac{1}{2}e^{-\frac{i\pi}{6}} & 0 & -\frac{1}{2}ie^{-\frac{i\pi}{6}} \\ 0 & 0 & -\frac{e^{-\frac{i\pi}{6}}}{2\sqrt{2}} & 0 & -\frac{1}{4}ie^{-\frac{i\pi}{6}} & 0 & 0 \end{pmatrix}$$

- We indeed have a restricted TQFT. Moreover, consider two torus bundles, determined by two matrices

$$A = \begin{pmatrix} 1 & 21 \\ 21 & 442 \end{pmatrix}, \quad B = \begin{pmatrix} 106 & 189 \\ 189 & 337 \end{pmatrix}$$

- We can easily calculate that we have

$$1 = \beta \circ (\rho(A) \otimes \mathbf{1}) \circ \gamma \neq \beta \circ (\rho(B) \otimes \mathbf{1}) \circ \gamma = 2$$

- We can distinguish between two torus bundles that are not distinguished by RT invariants.

- Along these lines, we wish to introduce the unit  $\varepsilon$  and the counit  $\eta$  (resulting in  $\beta\epsilon\mathbf{MCG}$  category, where **lens spaces** play a role), and finally multiplications and comultiplications (resulting in a  $\mu\epsilon\mathbf{MCG}$  category).
- In the former case, one can show that the additional relations are

$$\begin{aligned}\beta \circ (\mathbf{1} \otimes \varepsilon) &= \eta \\ D_a \circ \varepsilon &= \varepsilon\end{aligned}$$

- The way we construct those restricted TQFT's is similar to a dimensiona reduction from three to two dimensions, but with one important difference. We do not assume that our spacetime is of the form  $\Sigma \times S^1$ , but rather we have a nontrivial a la Hopf fibration.

**Thank you for your attention!**